

The Tadpole Conjecture in the Strict Asymptotic Regime

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Based on [arXiv:2204.05331]

with M. Graña, T. Grimm, D. van de Heisteeg, E. Plauschinn



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The Tadpole Conjecture in the Strict Asymptotic Regime

1. F-theory con Calabi-Yau four-folds and the Tadpole Conjecture
2. Asymptotic Hodge Theory \longrightarrow Strict Asymptotic Limit and Key Results
3. Moduli stabilization \longrightarrow General Results and Tadpole Scaling
4. Outlook & Open Questions

F-theory Compactifications

- Consider F-theory on a Calabi-Yau fourfold with fluxes

Review: [\[Denef '08\]](#)

Effective action:

[\[Grimm '10\]](#)

[\[Haack, Louis '21\]](#)

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- Restrict to complex structure sector: $J \wedge G_4 = 0$

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- Hodge decomposition and Hodge star: $H_{\text{prim}}^4 = H^{4,0} \oplus H^{3,1} \oplus H_{\text{prim}}^{2,2} \oplus H^{1,3} \oplus H^{0,4}$

$$\star v^{p,q} = i^{p-q} v^{p,q}$$

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[Bena, Blaback, Graña, Lüst '20]

Tadpole Conjecture: The flux contribution to the tadpole needed to stabilize a large number of moduli grows as

$$Q > \alpha n_{\text{stab}}$$

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Mariana's talk

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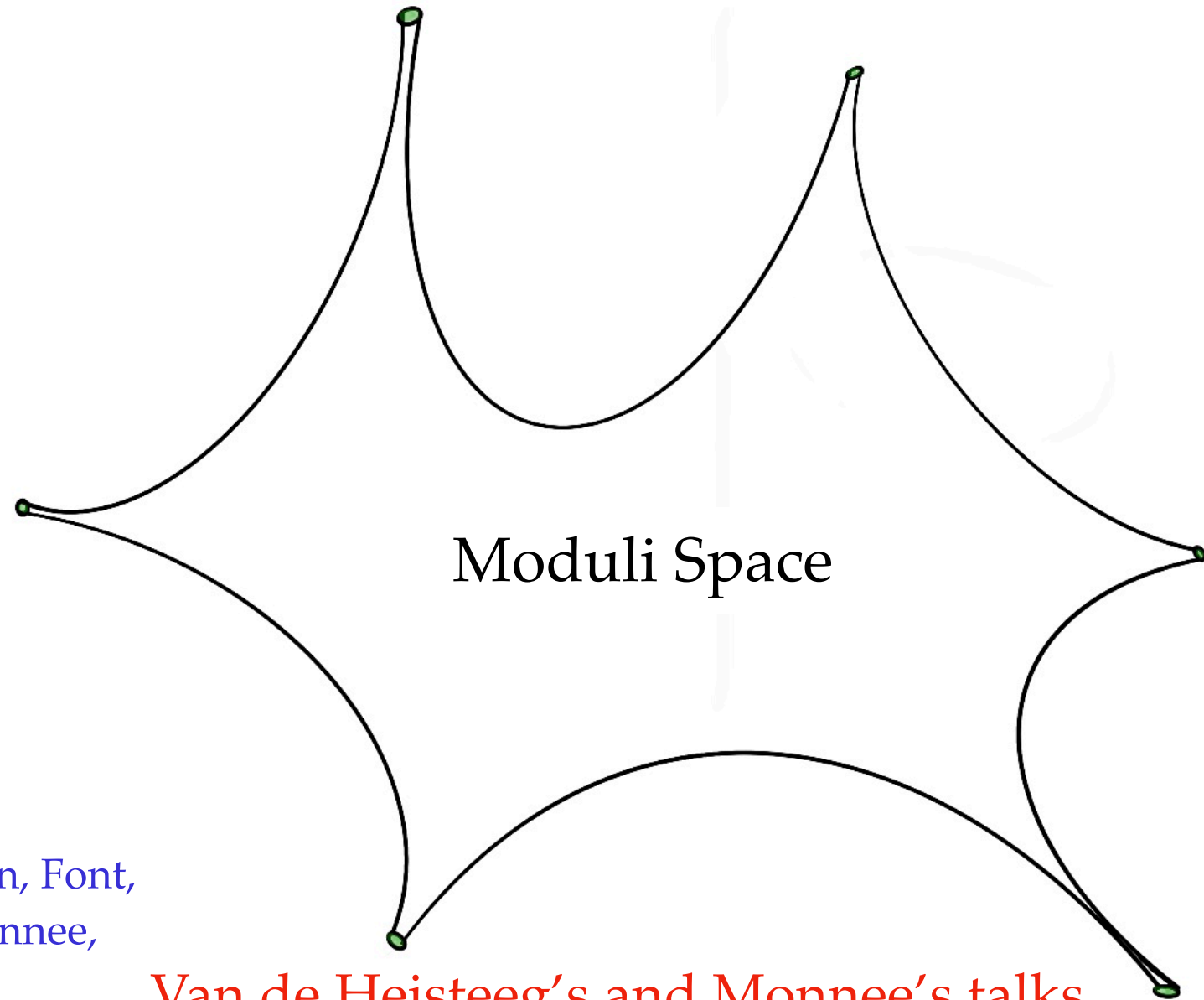
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GOAL: Prove this in the strict asymptotic region of moduli space

Asymptotic Hodge Theory

-Asymptotic limits-

[Griffiths, Deligne, Schmid,
Cattani, Kaplan...]



[Grimm, Palti, Valenzuela, Li, Bastian,
Castellano, Calderón-Infante, Corvilain, Font,
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Van de Heisteeg's and Monnee's talks

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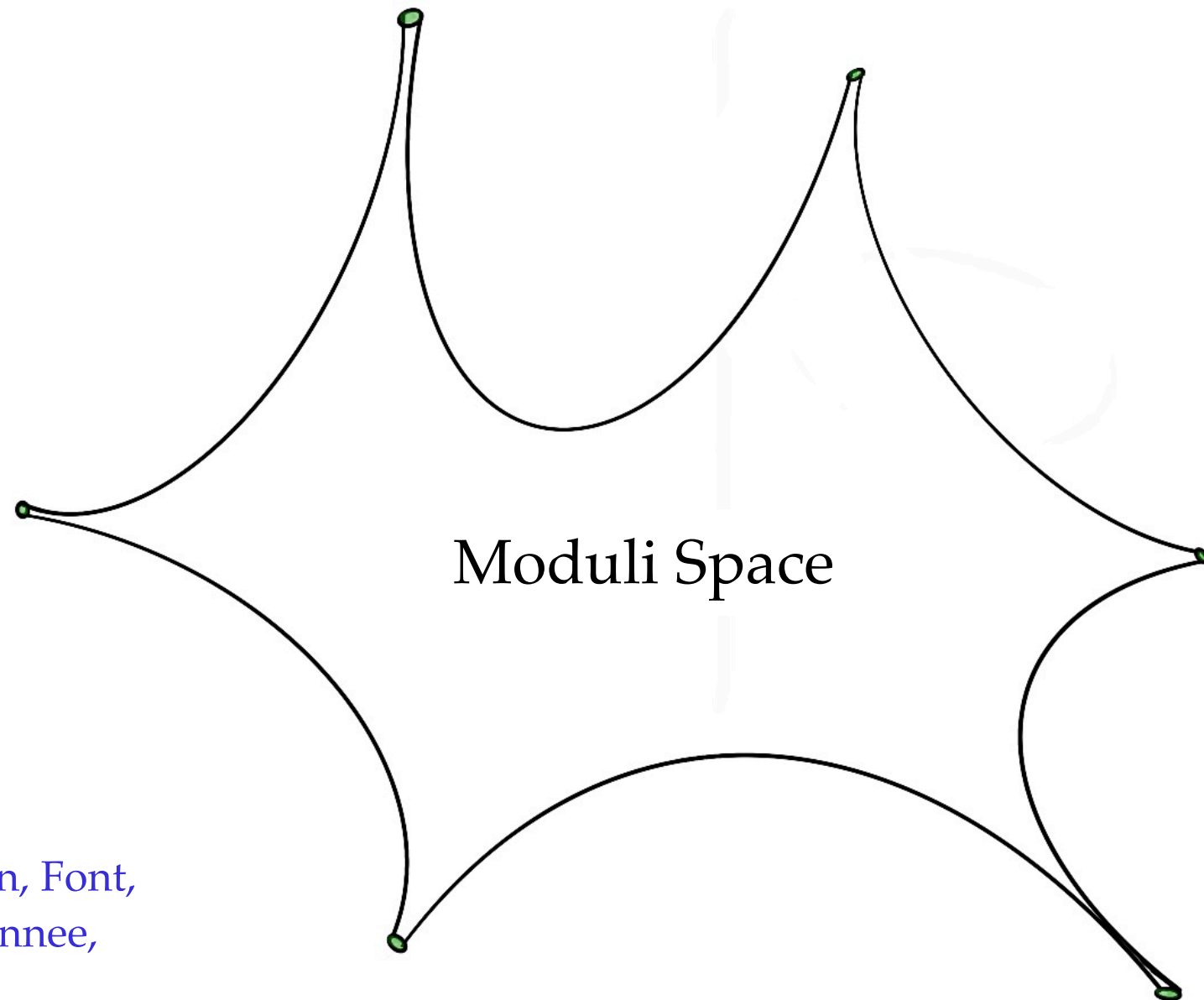
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with coordinates

$$t^i = \phi^i + is^i$$

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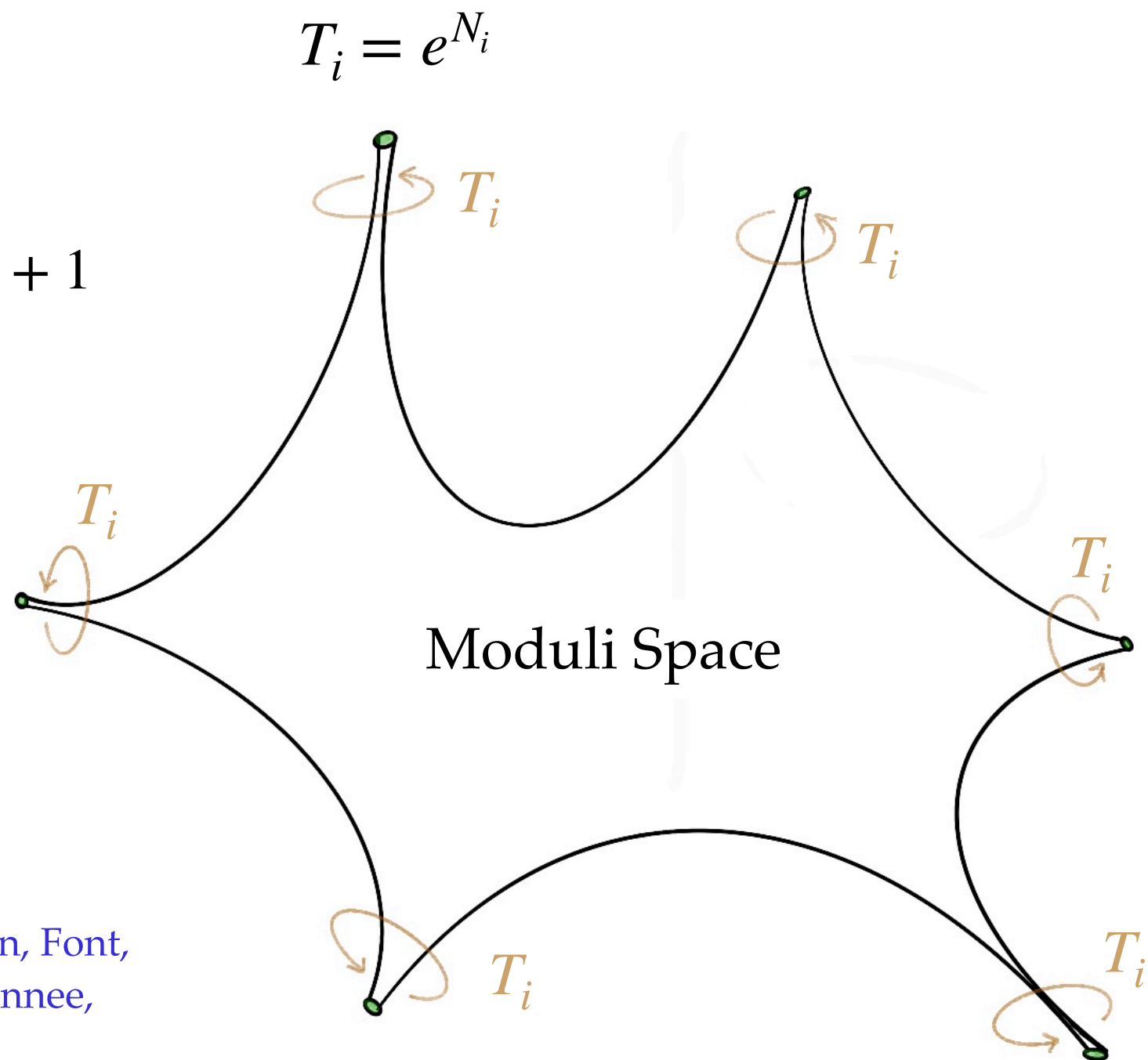
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$$\Pi(t^i + 1) = T_i \Pi(t^i)$$



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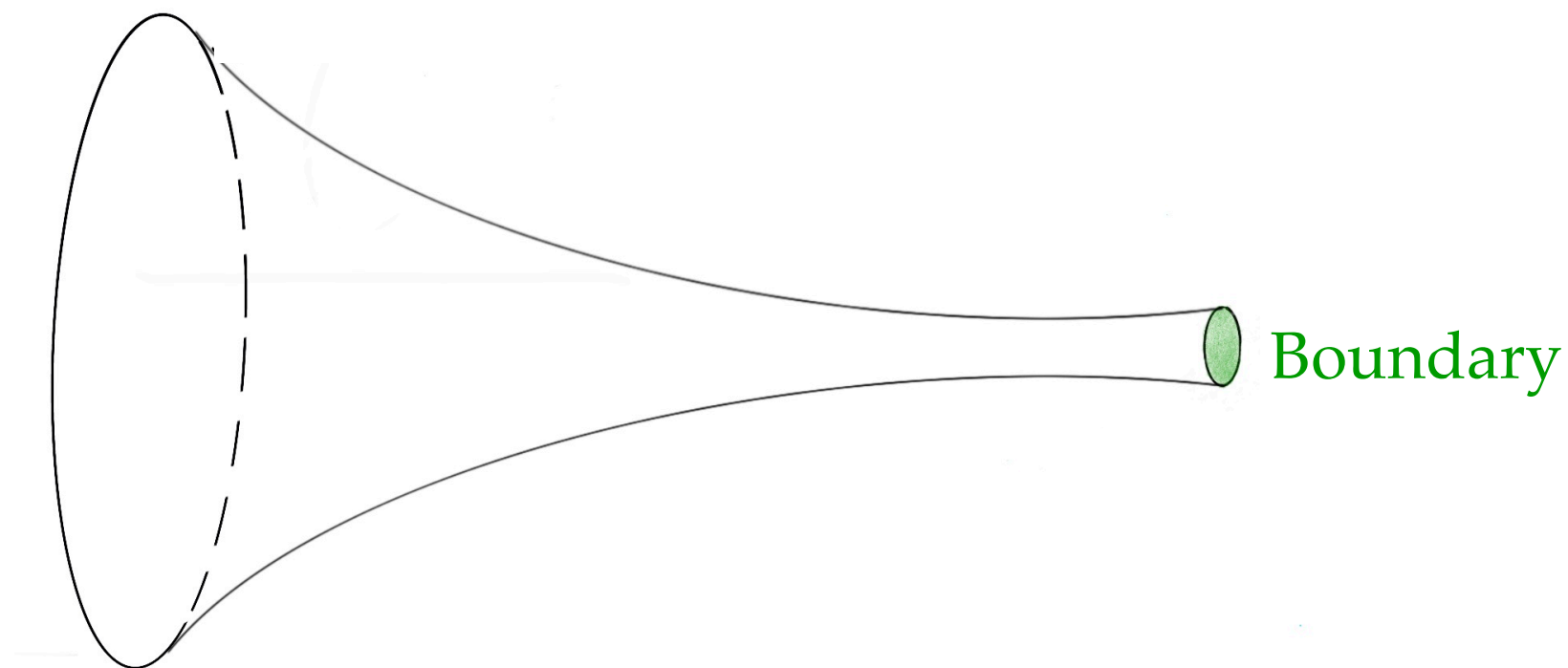
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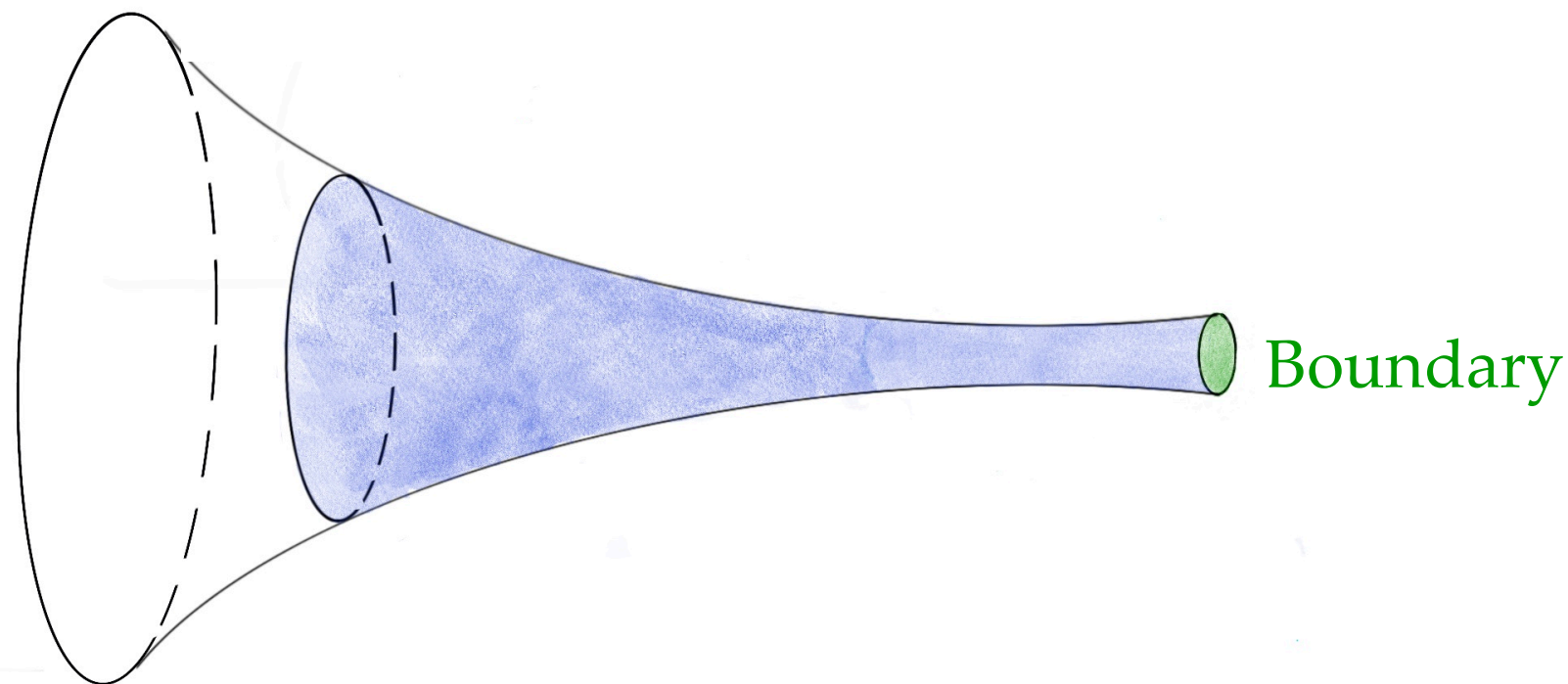
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Drop exponential corrections
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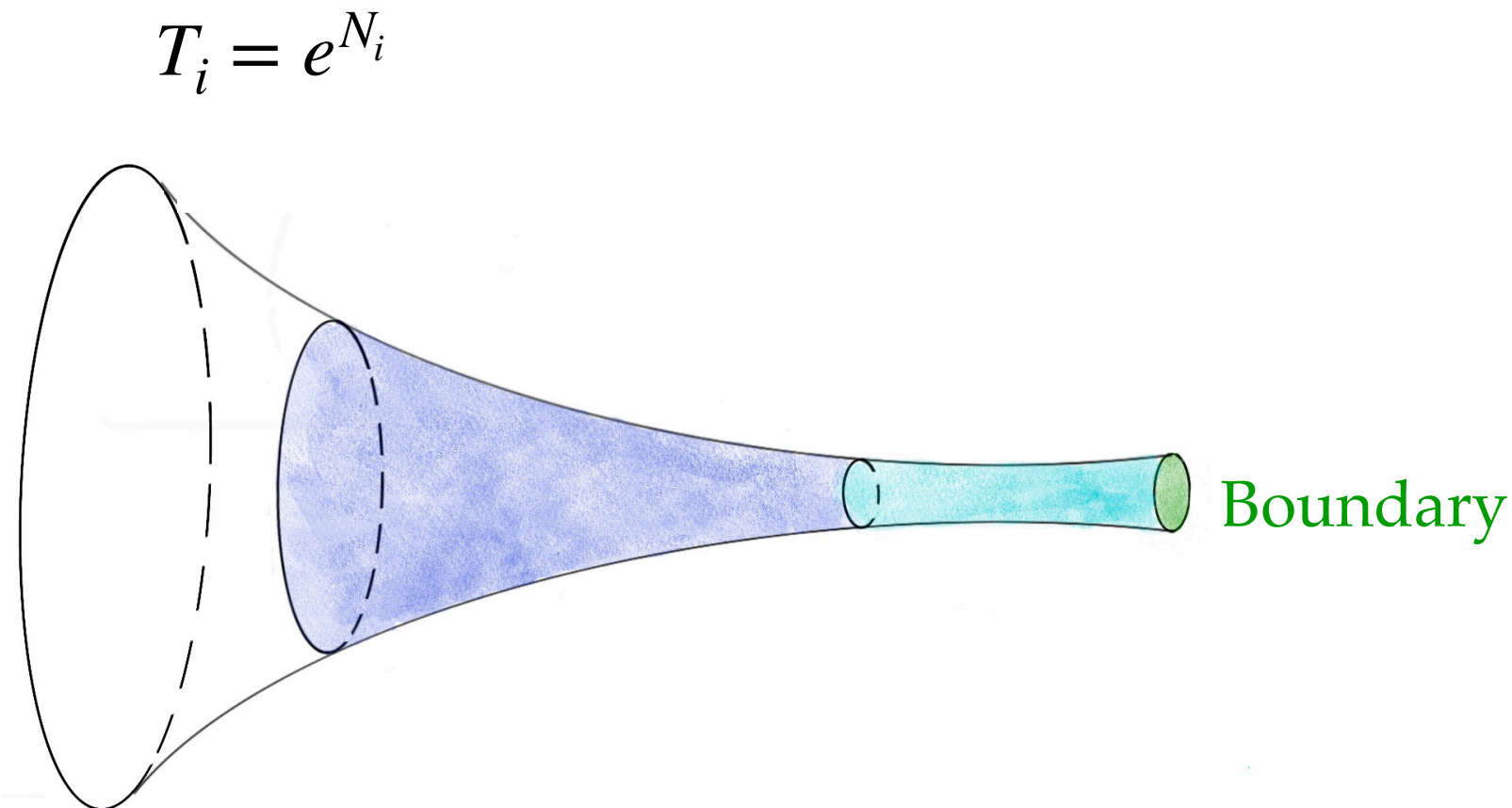
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- **Strict** Asymptotic Region \longrightarrow Introduce an **ordering** (drop polynomial corrections)

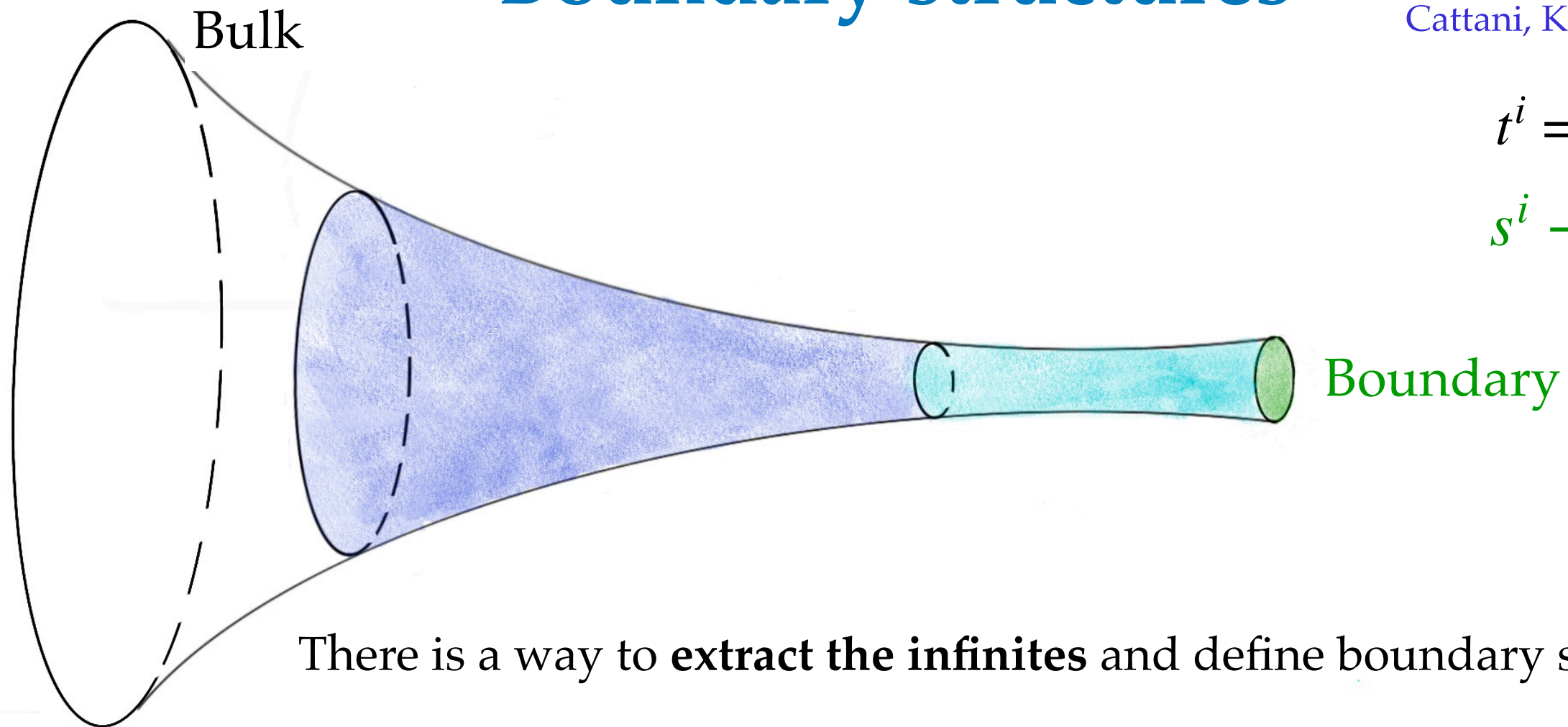
$$\frac{s^1}{s^2} > \gamma, \quad \frac{s^2}{s^3} > \gamma, \quad \dots, \quad \frac{s^{n-1}}{s^n} > \gamma, \quad s^n > \gamma \quad \gamma \gg 1$$



Asymptotic Hodge Theory

-Boundary structures-

[Griffiths, Deligne, Schmid, Cattani, Kaplan...]



$$t^i = \phi^i + i s^i$$

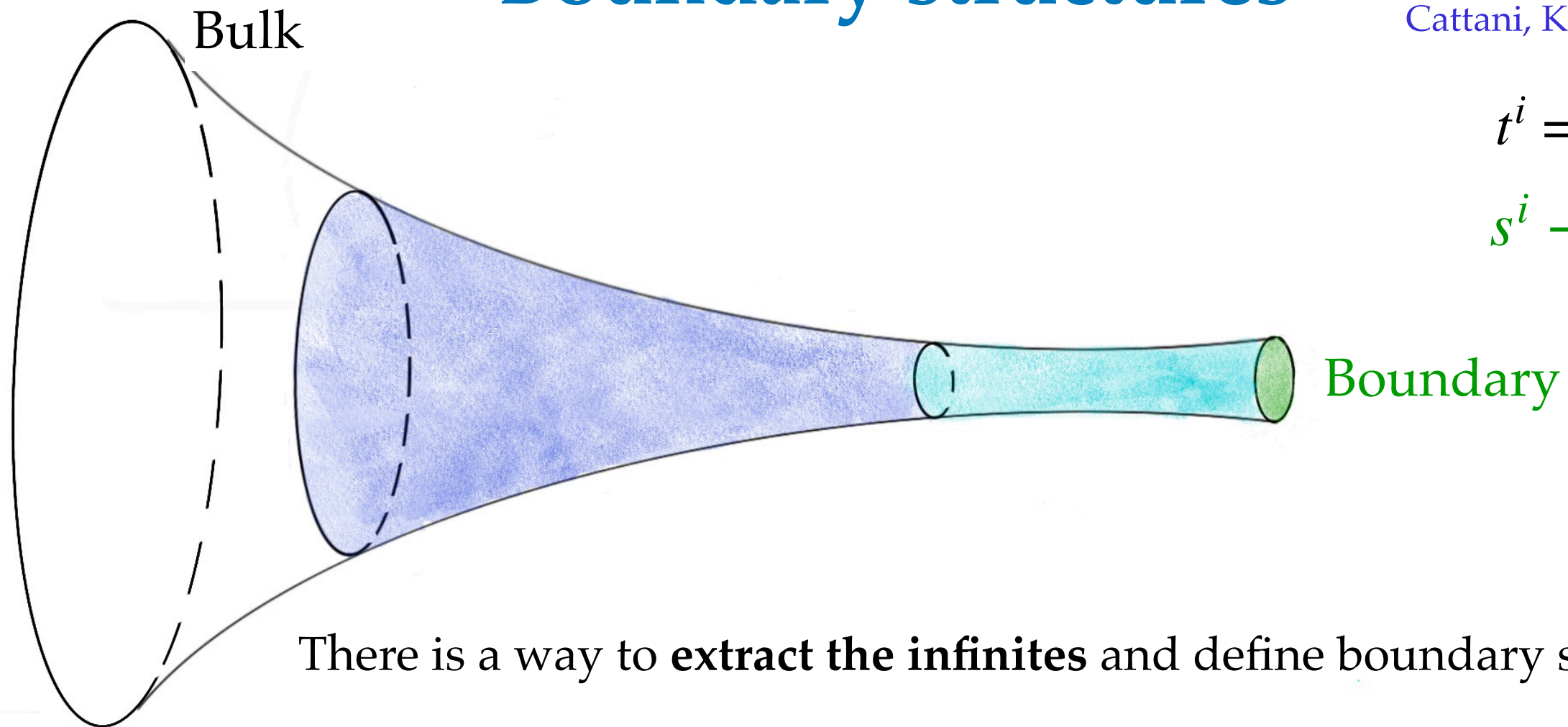
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There is a way to **extract the infinites** and define boundary structures

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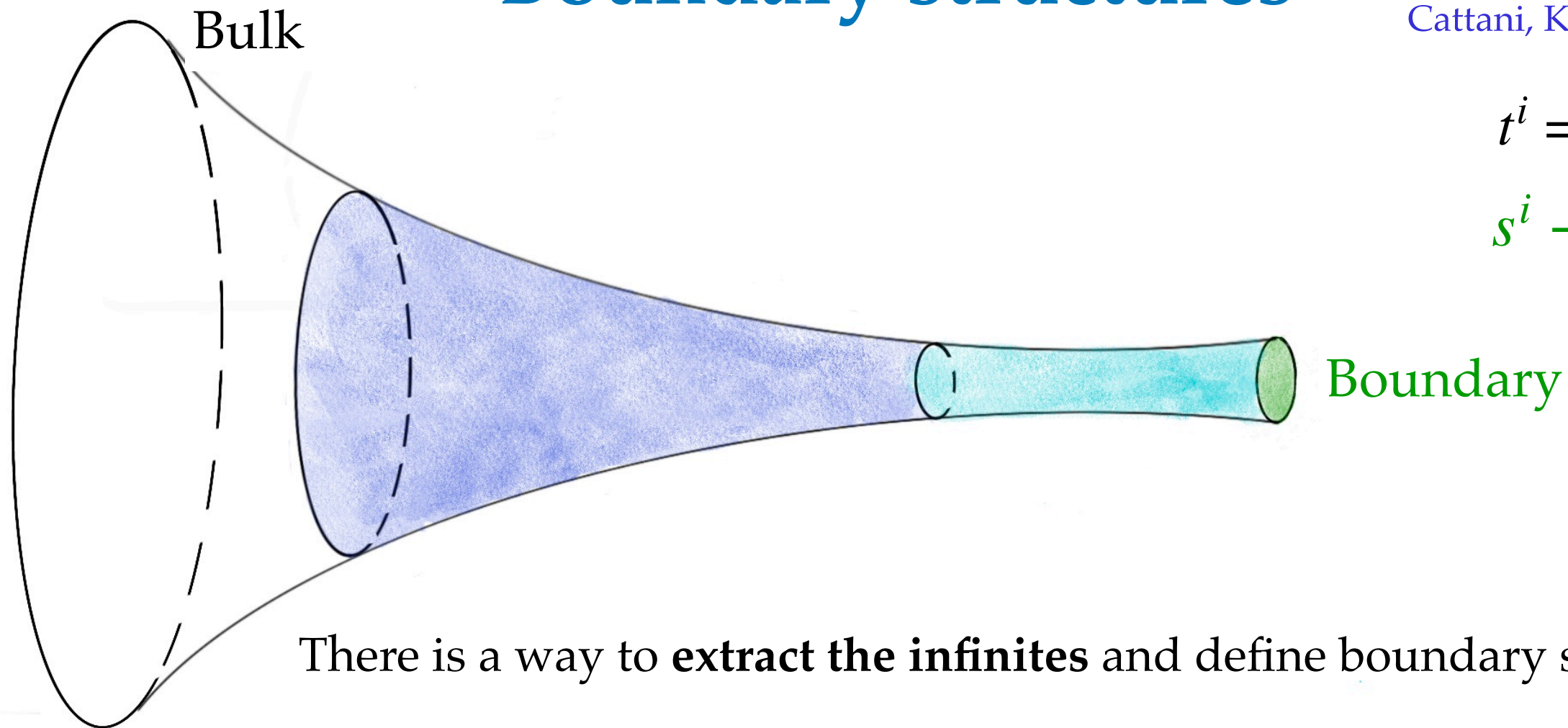
There is a way to **extract the infinites** and define boundary structures

- Boundary Hodge Decomposition: $H_{\text{prim}}^4(Y_4, \mathbb{C}) = H_{\infty}^{4,0} \oplus H_{\infty}^{3,1} \oplus H_{\infty}^{2,2} \oplus H_{\infty}^{1,3} \oplus H_{\infty}^{0,4}$

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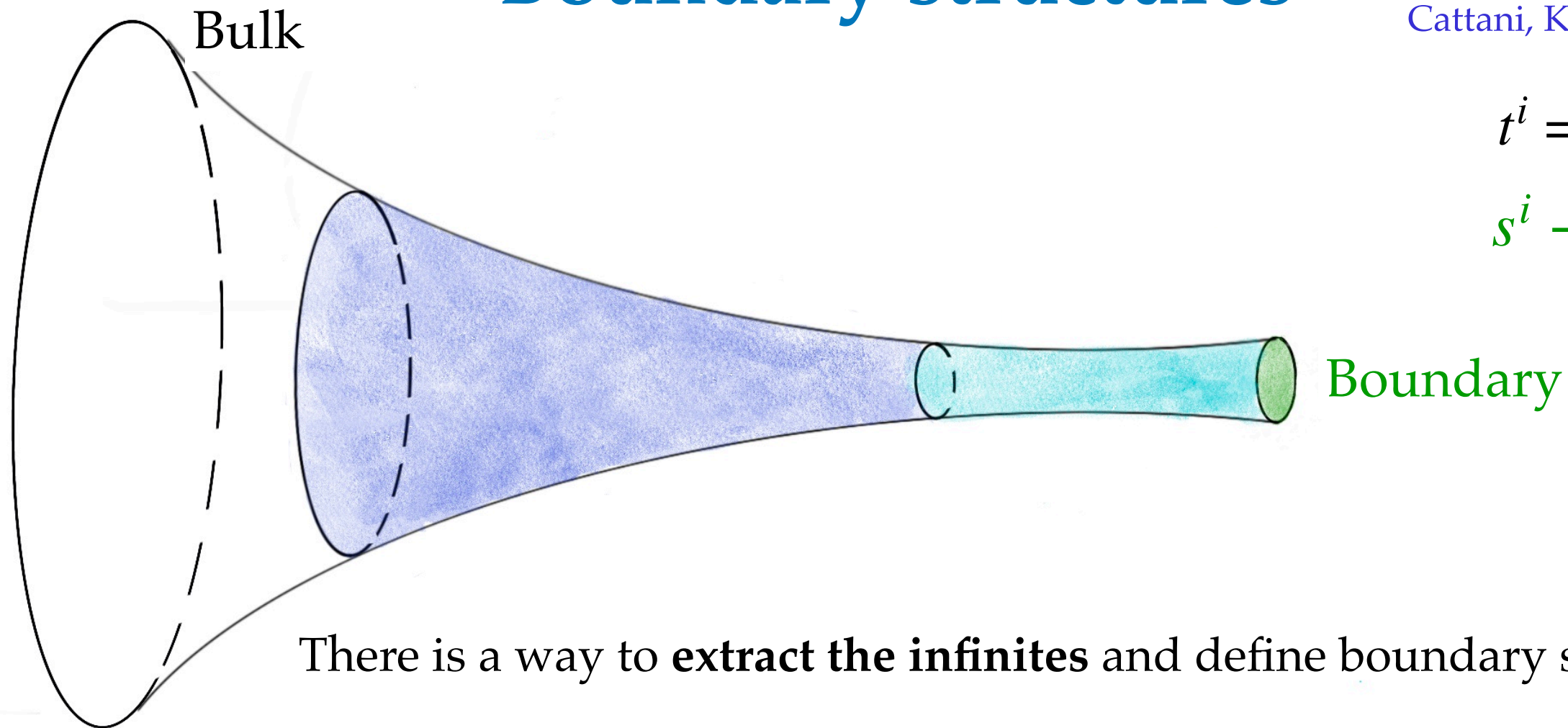


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- $\mathfrak{sl}(2)$ splitting: n **commuting** $\mathfrak{sl}(2)$ triplets: $\{N_i^-, N_i^+, N_i^0\}$

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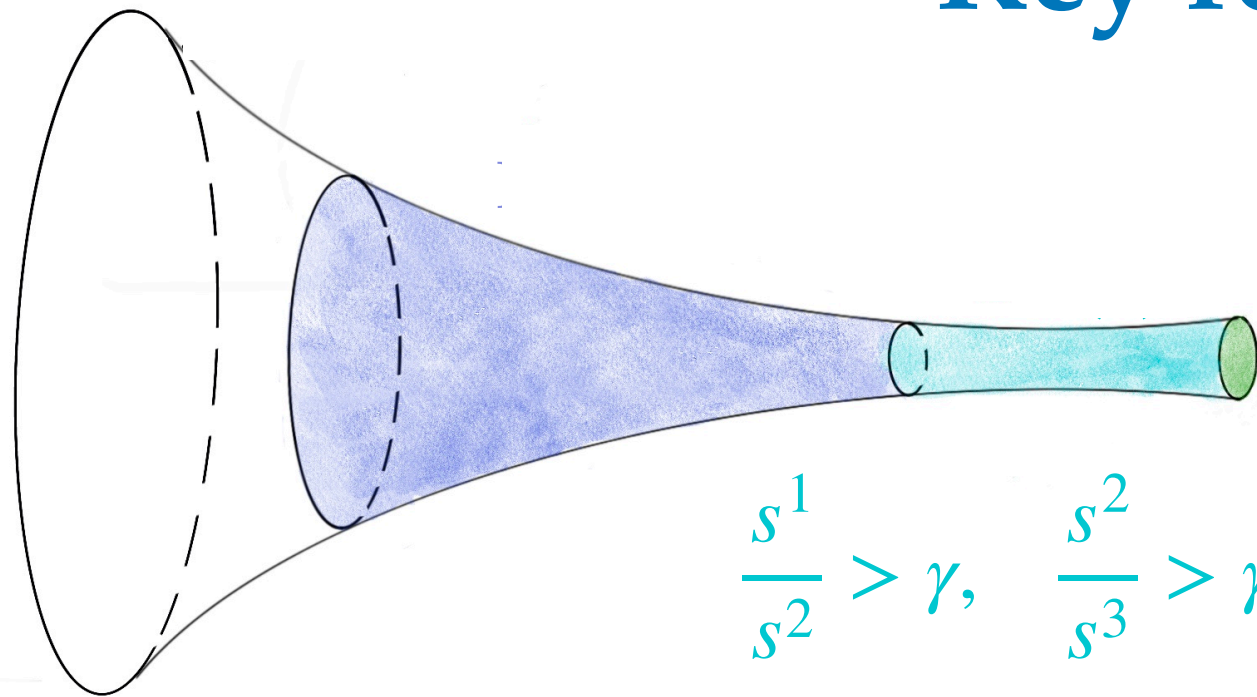


There is a way to **extract the infinities** and define boundary structures

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- Decompose fluxes into $\mathfrak{sl}(2)$ reps: $H_{\text{prim}}^4(Y_4, \mathbb{R}) = \bigoplus V_{\ell}$ $\ell = (\ell_1, \dots, \ell_n)$

Asymptotic Hodge Theory

-Key results-



$$\{N_i^-, N_i^+, N_i^0\} \quad H_\infty^{p, 4-p}$$

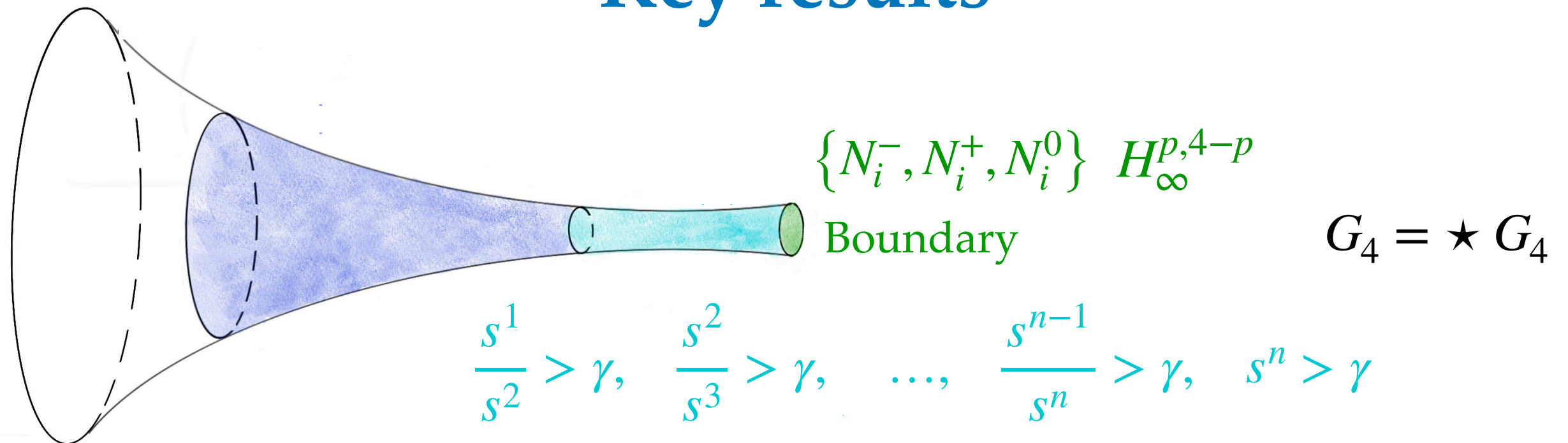
Boundary

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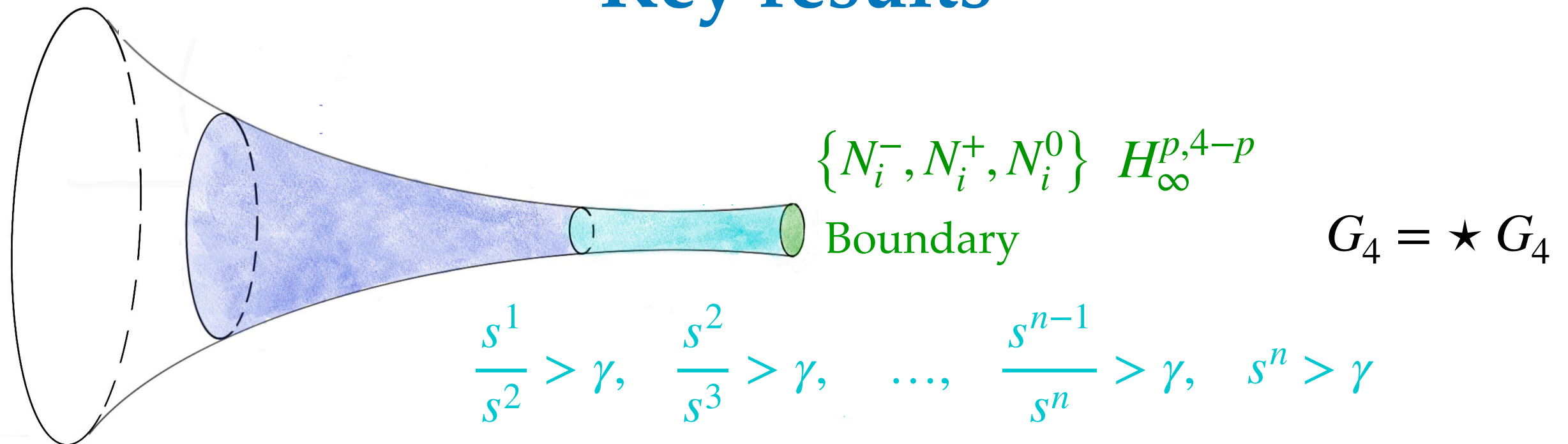
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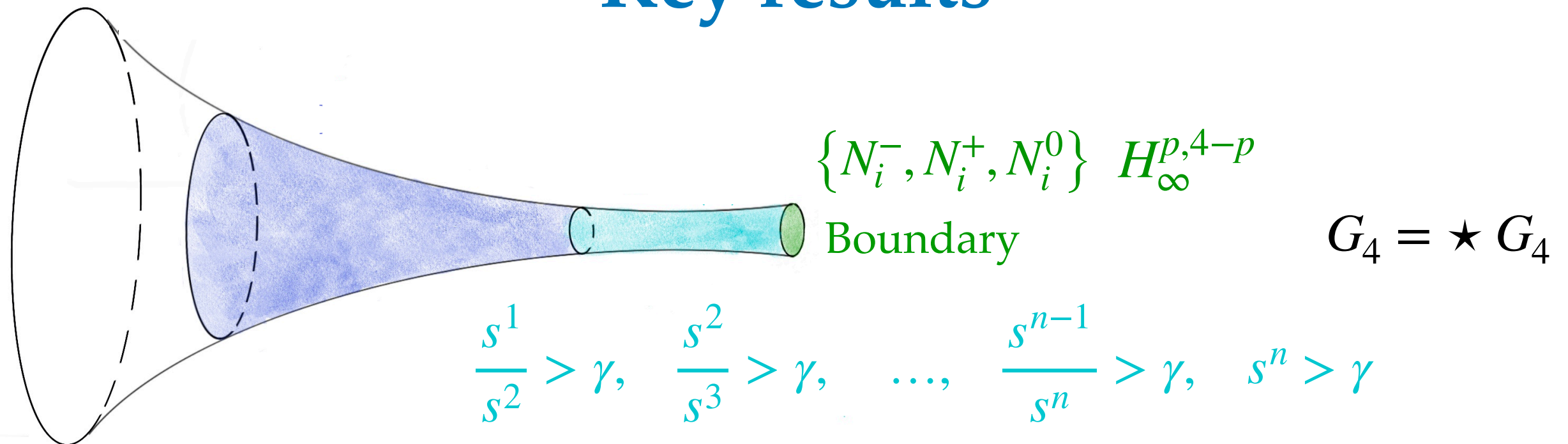
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 Only **four types**: P_0 P_{01_i} P_{02_i} $P_{01_i 2_j}$

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Only **four types**: $P_0 \quad P_{01_i} \quad P_{02_i} \quad P_{01_i 2_j}$
- Bdry and $\mathfrak{sl}(2)$ Hodge star: $\left(\begin{array}{l} \star_\infty : V_\ell \rightarrow V_{-\ell} \\ \star \longrightarrow \star_{\mathfrak{sl}(2)} \end{array} \right. \quad ||v_\ell||_{\mathfrak{sl}(2)}^2 = \left(\frac{s^1}{s^2}\right)^{\ell_1} \left(\frac{s^2}{s^3}\right)^{\ell_2} \dots \left(\frac{s^{n-1}}{s^n}\right)^{\ell_{n-1}} (s^n)^{\ell_n} ||v_\ell||_\infty^2$

Moduli stabilization: Generalities

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 - Combining many $\mathfrak{sl}(2)$ reps only imposes compatibility constraints between the fluxes, but never lowers the tadpole for a fixed modulus.
 - Tadpole-wise, the **most economic thing** is to turn on only **one $\mathfrak{sl}(2)$ per modulus**.

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At least linear scaling with (a large number of) stabilized moduli

Moduli stabilization and Tadpole

-Key results-

- For large number of stabilized moduli, most of them (i.e. at least $(n - 4)$ of them) have to be fixed with fluxes from these representations

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Outlook & Open Questions

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Exclude scaling of $||G_\ell||_\infty^2 \sim \frac{1}{n \gamma^{\sum_i \ell_i}}$ \longrightarrow Found numerical evidence

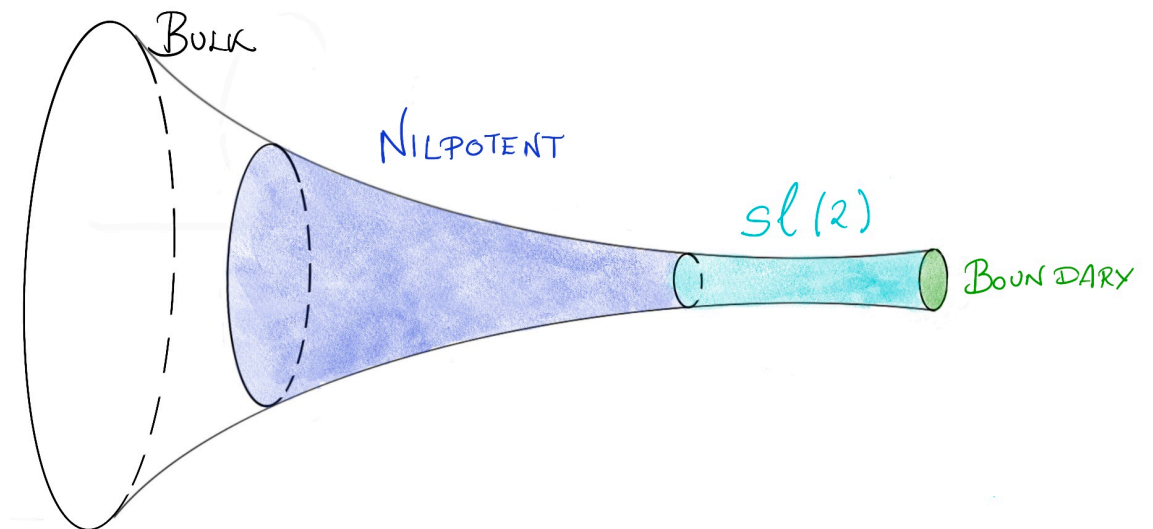
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- Extend to the **interior of moduli space**

- 1) Include polynomial corrections (strict asymptotic limit) \longrightarrow “Linear scenario”

[Grimm '20]

[Marchesano, Prieto, Wiesner '21]

[Palti, Tasinato, Ward '08]

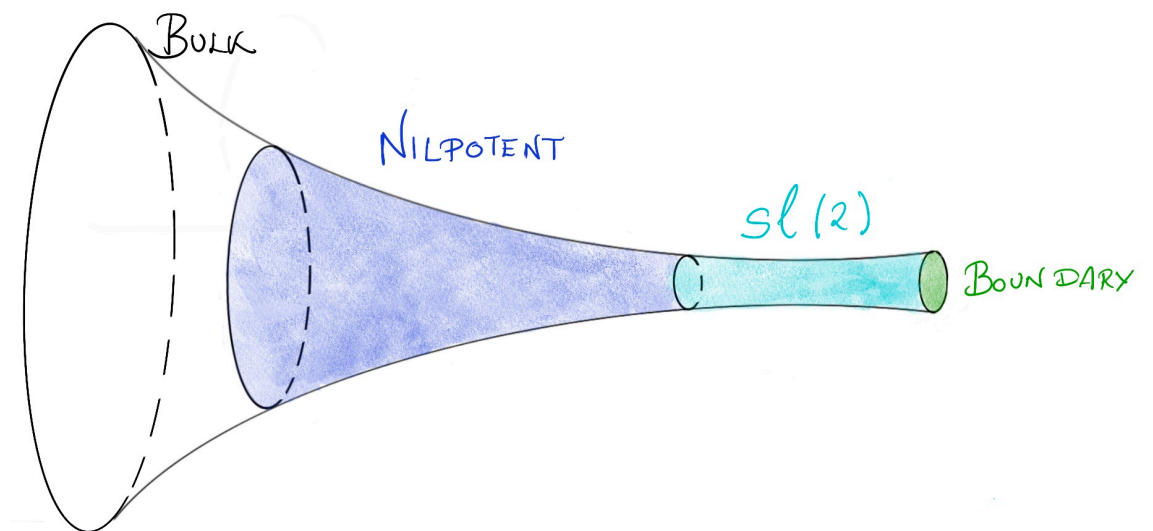
See also: [Plauschinn '22] [Lüst '22]

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- 2) Interior of moduli space \longrightarrow Hodge loci of $G_4^{2,2}$ fluxes is algebraic

[Bakker, Grimm, Schnell, Tsimmerman '21]

[Cattani, Deligne, Kaplan '95]

A word cloud featuring the phrase "Thank You" in numerous languages. The words are arranged in a circular pattern, with "thank you" in the center in large, bold, lowercase letters. Surrounding it are many other languages, including English ("thanks", "gracias", "merci", "thank you"), Spanish ("gracias", "muchas gracias", "muchas gracias"), French ("merci", "merci beaucoup", "merci"), German ("danke", "danke"), Italian ("grazie", "grazie"), Japanese ("arigato", "arigato gozaimasu", "arigato"), Chinese ("谢谢", "谢谢"), Korean ("고맙습니다", "고맙습니다"), Hindi ("धन्यवाद", "धन्यवाद"), and many others. The words are in various colors and sizes, creating a vibrant and multicultural visual.

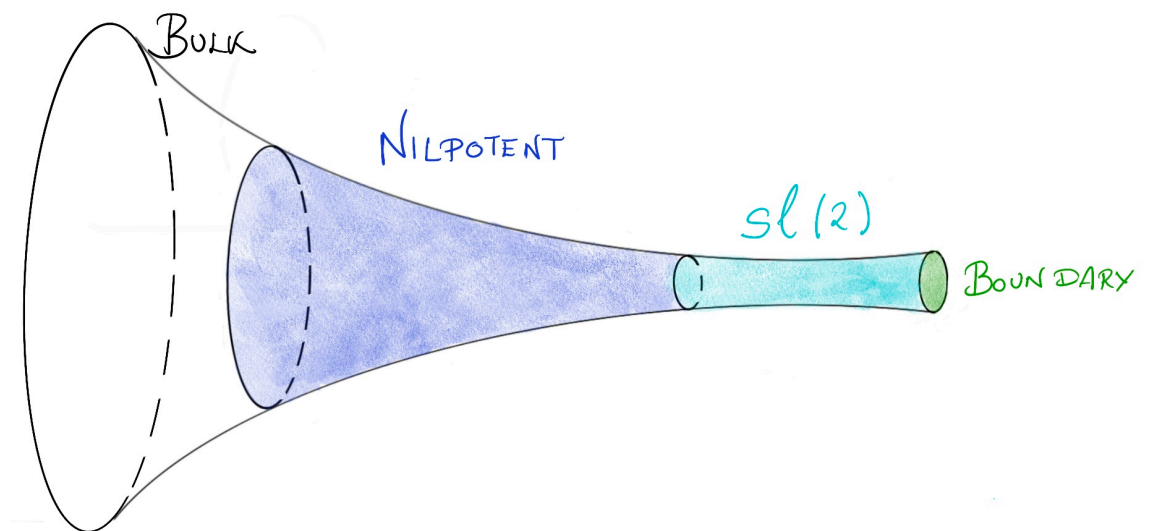
Backup slides

Charge quantization and extensions

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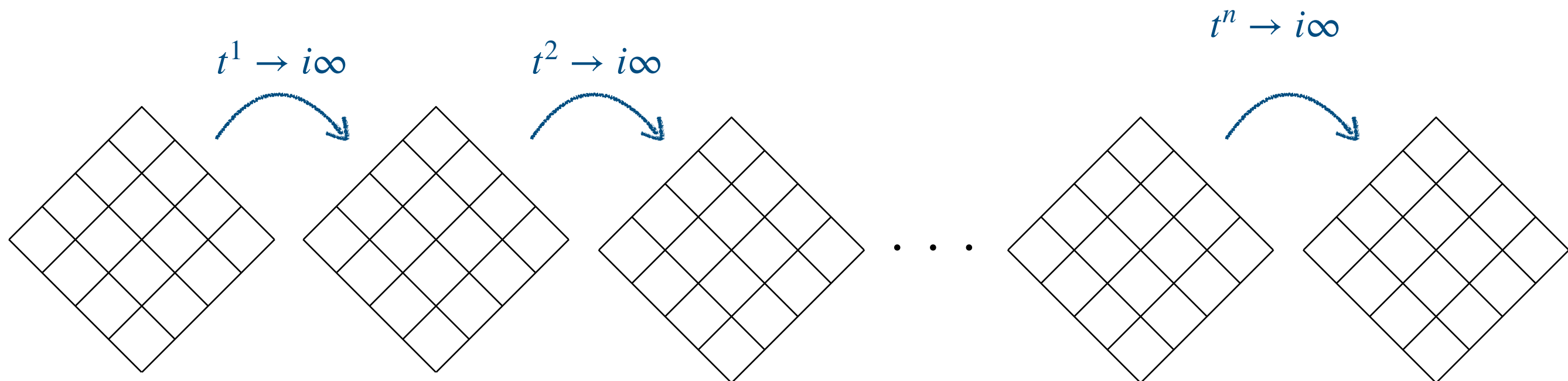
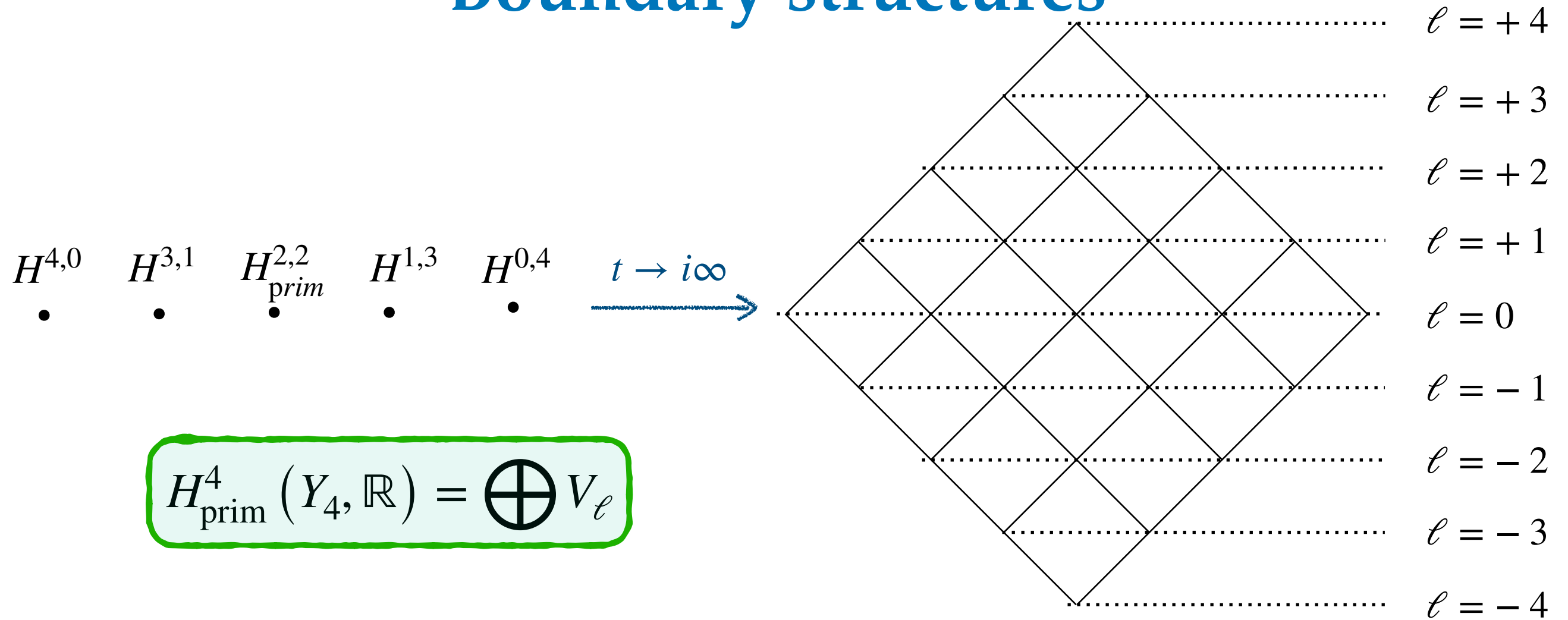
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Asymptotic Hodge Theory

-Boundary structures-



Asymptotic Hodge Theory

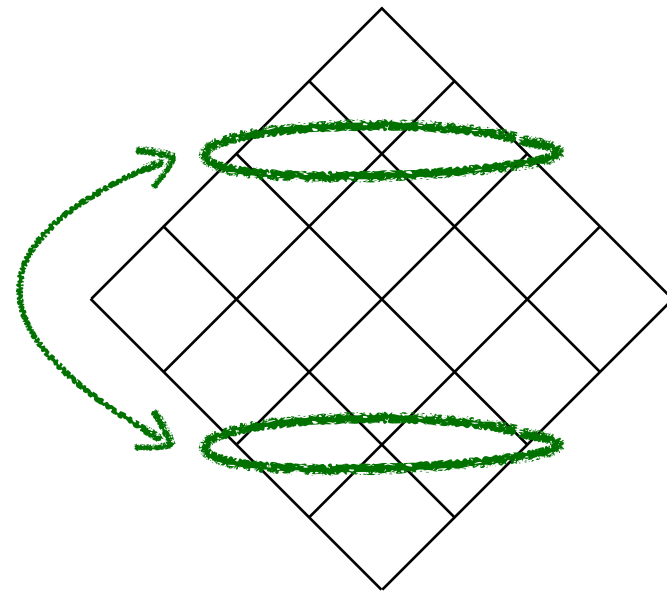
-Hodge star close to the boundary-

- The boundary Hodge decomposition naturally includes a boundary Hodge star operator:

$$\star_{\infty} : V_{\ell} \rightarrow V_{-\ell}$$

$$\langle v_{\ell}, v_{\ell'} \rangle = 0 \quad \text{for } \ell \neq -\ell'$$

$$\langle v_{\ell}, \star_{\infty} v_{\ell'} \rangle = 0 \quad \text{for } \ell \neq \ell'$$



- Allows to express the Hodge star in the strict asymptotic limit: $\star \longrightarrow \star_{\text{sl}(2)}$

Vanishing axions: $\star_{\text{sl}(2)} v_{\ell} = \left(\frac{s^1}{s^2} \right)^{\ell_1} \left(\frac{s^2}{s^3} \right)^{\ell_2} \cdots \left(\frac{s^{n-1}}{s^n} \right)^{\ell_{n-1}} (s^n)^{\ell_n} \star_{\infty} v_l$

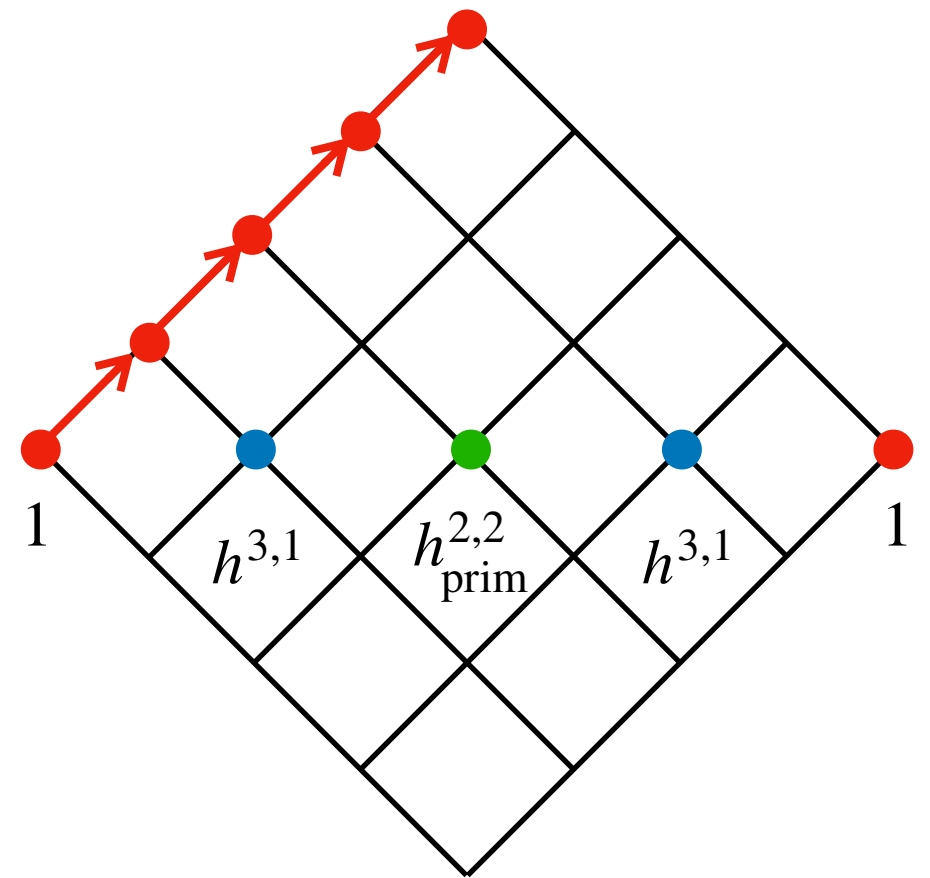
Non-vanishing axions \longrightarrow Mixing with lower subspaces via $e^{\phi^i N_i^-} v_l$

Asymptotic Hodge Theory

-Highest weight spaces-

- Elements in $H_{\text{prim}}^4(Y_4, \mathbb{R})$ arrange into irreps of the n commuting $\mathfrak{sl}(2) \longrightarrow$
Generated by applying N_i^- to **highest weight states**
- What are the possible highest weight states?

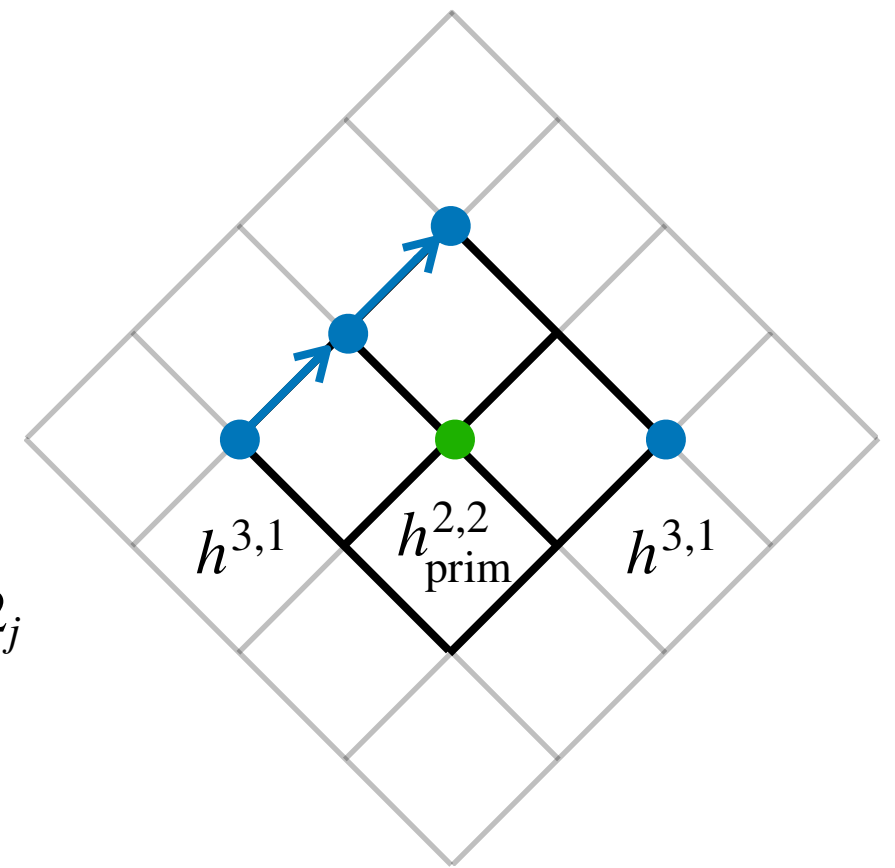
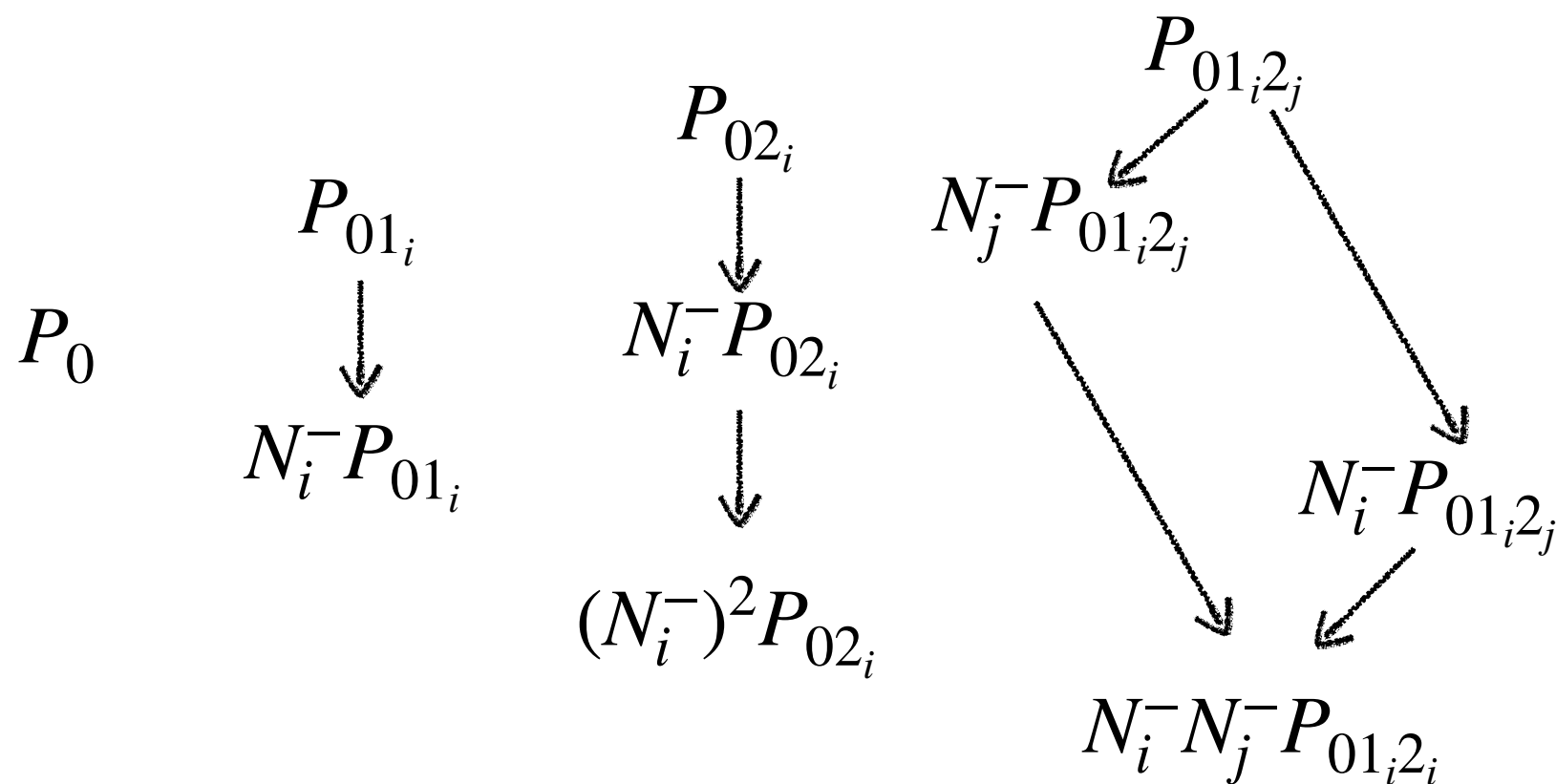
Only **ONE**, corresponding to the **(4,0)-form**, can move along the exterior line of the diagram



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$\sim h^{3,1}$ copies of K3 surfaces

Moduli stabilization at work

$$\hat{G}_{-\ell} = \left(\frac{s^1}{s^2} \right)^{\ell_1} \cdots \left(\frac{s^{n-1}}{s^n} \right)^{\ell_{n-1}} (s^n)^{\ell_n} \star_{\infty} \hat{G}_{+\ell}$$

$$\hat{G}_4 = e^{\phi^i N_i^-} G_4$$

- Example: Flux from the $\mathfrak{sl}(2)$ rep generated by P_{02_i}

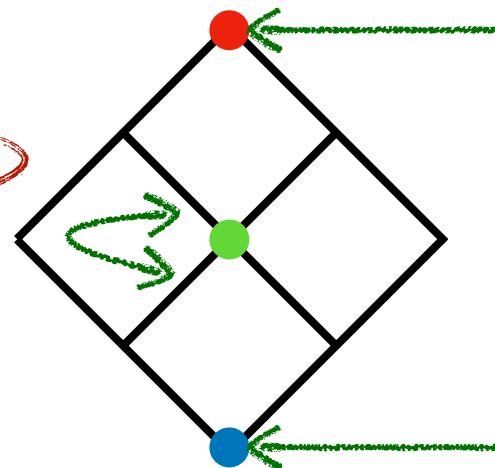
$$G_4 = G_{02} \textcolor{red}{v}_{02} + G_0 \textcolor{green}{v}_0 + G_{0-2} \textcolor{blue}{v}_{0-2} \longrightarrow Q = \frac{1}{2} \langle G_4, G_4 \rangle = G_{02} G_{0-2} - \frac{1}{2} G_0^2$$

$$\hat{G}_4 = G_{02} \textcolor{red}{v}_{02} + (G_0 - \phi G_{02}) \textcolor{green}{v}_0 + \left(G_{0-2} - \phi G_0 + \frac{1}{2} (\phi)^2 G_{02_i} \right) \textcolor{blue}{v}_{0-2}$$

- Self-duality conditions:

$$\frac{G_{02}}{2} (s)^2 = G_{0-2} - \phi G_0 + \frac{1}{2} (\phi)^2 G_{02} \longrightarrow s = \frac{\sqrt{2G_{0-2}G_{02} - G_0^2}}{G_{02}}$$

$$(G_0 - \phi G_{02}) = - (G_0 - \phi G_{02}) \longrightarrow \phi = \frac{G_0}{G_{02}}$$



sl(2) Hodge star

- Extend boundary Hodge star to the interior: $\star_\infty \longrightarrow \star_{\text{sl}(2)}$

$$\star_{\text{sl}(2)} = e^{+\phi^i N_i^-} \left[e^{-\frac{1}{2} \log(s^i) N_i^0} \star_\infty e^{+\frac{1}{2} \log(s^i) N_i^0} \right] e^{-\phi^i N_i^-}$$

$$H_{\text{prim}}^4(Y_4, \mathbb{C}) = H_{\text{sl}(2)}^{4,0} \oplus H_{\text{sl}(2)}^{3,1} \oplus H_{\text{sl}(2)}^{2,2} \oplus H_{\text{sl}(2)}^{1,3} \oplus H_{\text{sl}(2)}^{0,4}$$

$$H_{\text{sl}(2)}^{p,q} = e^{\phi^i N_i^-} e^{-\frac{1}{2} \log(s^i) N_i^0} H_\infty^{p,q}$$